

The Crystal Structure of S-Methylisothiurea Sulphate $[\text{CH}_3\text{SC}(\text{NH}_2)_2]_2^+\text{SO}_4^{2-}$

By C. H. STAM

Laboratory for Crystallography, University of Amsterdam, The Netherlands

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Crystals of S-methylisothiurea sulphate are orthorhombic with space group *Pcan*. The unit-cell dimensions are:

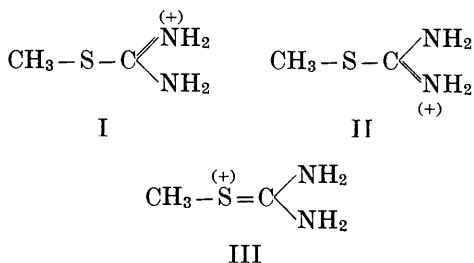
$$a = 8.38, b = 11.32, c = 12.60 \text{ \AA}.$$

The structure was refined by least squares using three-dimensional X-ray data obtained with Cu $K\alpha$ radiation. The final *R*-factor for 882 observed and 324 non-observed reflections was 15.2%.

The $\text{SC}(\text{NH}_2)_2$ part of the molecule is planar. The two C-S bonds have lengths of 1.79 and 1.74 \AA. The C-S-C angle is 104° . The SO_4 ion does not deviate significantly from a regular tetrahedron; the average S-O distance is 1.463 \AA.

Introduction

Simple valence-bond considerations predict resonance in the S-methylisothiurea ion between configurations I, II and III:



This would lead to one essentially single C-S bond and one of partial double-bond character. One would expect III to have the largest contribution since here the positive charge is on the less electronegative atom.

An accurate structure determination has been reported for thiurea by Kunchur & Truter (1958), and it is interesting to compare their results with those for S-methylthiurea sulphate. An additional point of interest is the SO_4 -ion for which few accurate data are available.

Unit-cell data. Intensities

The crystals are orthorhombic. The following unit-cell dimensions were obtained from Straumanis-type oscillation diagrams:

$$a = 8.38 \pm 0.01, b = 11.32 \pm 0.01, c = 12.60 \pm 0.01 \text{ \AA}.$$

By flotation the density was found to be 1.54 g.cm^{-3} , agreeing with the calculated value for $Z=4$.

Systematic absences lead to the space group *Pcan*(*Pbcn*).

The crystals at our disposal were for the greater part needles in the [001] direction with the form {110} well developed and small faces of the forms {010},

{111}, {121} and {012}. A few crystals had a more isometric habit. Measurements on several crystals with the optical two-circle goniometer yielded $a:b:c = 0.740:1:1.117$, as compared with $0.740:1:1.113$ from the above X-ray data. The crystals showed good cleavage along (110). Weissenberg diagrams of the layers zero through four about the *a*-axis and zero through seven about the *b*-axis were taken with Cu $K\alpha$ radiation using the multiple-film technique. The intensities were measured visually by comparison with an intensity scale. Absorption corrections were not applied although absorption was not negligible, the cross sections of the crystals having dimensions of about 0.15 and 0.20 mm. for the *a*- and *b*-axis respectively, and μ being 52.2 cm^{-1} .

The total number of independent reflections measured was 882. The number that could have been recorded on our films is 1206, the number within the reflection sphere being 1350 (systematic extinctions excluded). The 324 reflections that escaped observation were included in the least-squares refinement with values corresponding to one half the minimum observable intensities.

Structure determination

The structure was solved by means of Patterson and electron-density projections. There is



in the asymmetric unit and thus the sulphur atom of sulphate ion has to be at a special fourfold position. There are three fourfold positions in the space group *Pcan*, of which only $4c: x, 0, \frac{1}{4}; \bar{x}, 0, \frac{3}{4}; \frac{1}{2}+x, \frac{1}{2}, \frac{1}{4}; \frac{1}{2}-x, \frac{1}{2}, \frac{3}{4}$; need be considered; the other two, being at centres of symmetry, are ruled out by the tetrahedral configuration of the sulphate ion. In the $0kl$ projection the position of the sulphate sulphur atom is thus fixed; those of the other sulphur and the oxygen atoms in the asymmetric unit could be approximately

found from the Patterson projection. Starting from these sulphur and oxygen coordinates it was possible to locate the other atoms by means of successive electron-density projections. A similar procedure was used for the $h0l$ projection. These two projections were refined to R -factors of 21 and 22% respectively. The third projection along the c -axis was used to check the three-dimensional structure deduced from the a - and b -axis projections.

A three-dimensional least-squares refinement was then started. After three cycles with individual isotropic temperature factors the R factor had dropped from 26.7 to 19.6%. As it was clear from the projections that there was marked anisotropy in the electron distribution of the sulphur atoms, anisotropic temperature parameters were introduced for these atoms after the third cycle. After the fourth cycle R dropped to 16.1%. A number of low-order reflections showed rather poor agreement, caused partly by extinction. In order to examine the effect of this on the refinement, only reflections with $\sin \theta > 0.5$ were used in the next two cycles; R now dropped to 14.4%. When the low-order reflections after some minor improvements were included again, R for all reflections was 17.6%. Two more cycles were calculated, the final R being 15.2%. The final coordinates do not differ significantly from those of the refinement with the low orders omitted (largest difference 0.015 Å in the y of C_1 ; see Table 2). The final temperature factors, however, are consistently lower, but because of neglect of the absorption correction and omission of the hydrogen atoms the temperature parameters are not considered to be very reliable. A list of observed and calculated structure factors is given in Table 1. The scattering factors of Berghuis *et al.* (1955) for C, N and O, and those of Tomiie & Stam (1958) for S were used.

Results

In Table 2 the final coordinates are listed together with the calculated standard deviations. Bond distances are given in Fig. 1(a) and Table 3, bond angles in Fig. 1(b) and Table 4. Table 5 lists the temperature parameters. The principal vibration directions of the two sulphur atoms are given in Table 6. In Fig. 2 and 3 projections of the structure along the a - and c -axes are shown with the hydrogen bonds indicated by broken lines. The probable course of the (110)

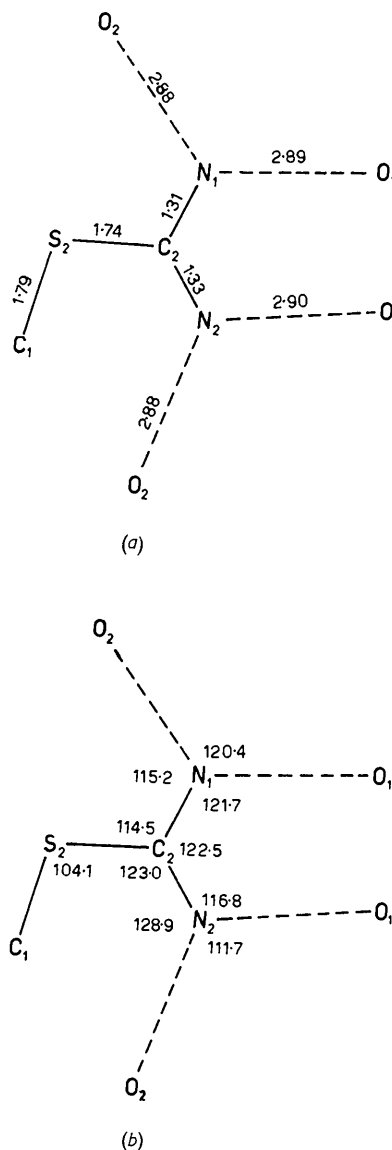


Fig. 1. Bond length and angles in the S-methylisothiourea ion.

cleavage plane is indicated by the chain line in Fig. 3. In Fig. 2 the interionic distances shorter than 3.50 Å are indicated.

Table 2. Final coordinates and calculated standard deviations in Å

Between brackets the coordinates resulting from the refinement with the low orders omitted

	x	$\sigma(x)$	y	$\sigma(y)$	z	$\sigma(z)$
S_1	3.5972	(3.5981)	0.0022	0.0000	(0.0000)	0.0000
S_2	7.3531	(7.3500)	0.0023	2.6644	(2.6673)	0.0022
O_1	4.4412	(4.4445)	0.0052	0.9698	(0.9679)	0.0055
O_2	2.7519	(2.7526)	0.0050	0.6944	(0.6961)	0.0051
N_1	6.2121	(6.2114)	0.0063	2.1957	(2.1957)	0.0066
N_2	5.9691	(5.9686)	0.0059	0.4592	(0.4660)	0.0063
C_1	7.9328	(7.9338)	0.0090	1.5174	(1.5319)	0.0088
C_2	6.4084	(6.4061)	0.0066	1.6739	(1.6732)	0.0072

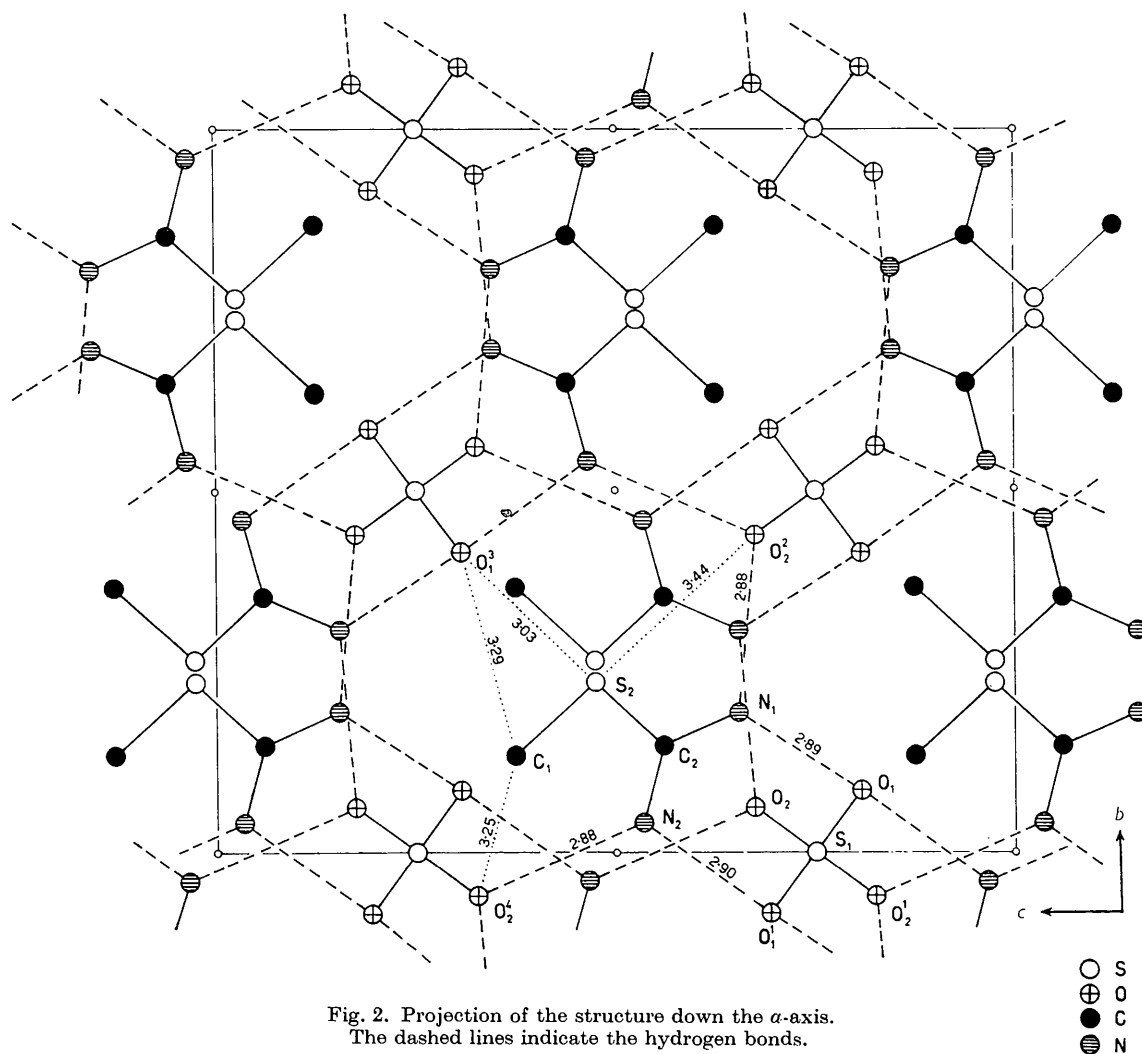


Fig. 2. Projection of the structure down the a -axis. The dashed lines indicate the hydrogen bonds.

Table 3. Bond distances and calculated standard deviations

	Distance	σ
C_1-S_2	1.790 Å	0.009 Å
C_2-S_2	1.743	0.007
C_2-N_1	1.308	0.009
C_2-N_2	1.333	0.009
S_1-O_1	1.470	0.006
S_1-O_2	1.456	0.006
$N_1 \cdots O_1$	2.889	0.008
$N_1 \cdots O_2$	2.875	0.008
$N_1 \cdots O_1^1$	2.903	0.008
$N_2 \cdots O_2^4$	2.878	0.008

Discussion

The thiourea part of the molecule comprising S_2 , C_2 , N_1 and N_2 is practically planar. The best plane through these atoms, calculated with the method of Schomaker *et al.* (1959), is:

$$0.8435x - 0.4151y - 0.3409z = 2.8361.$$

Table 4. Bond angles and calculated standard deviations

	Angle	σ
$C_1-S_2-C_2$	104.1°	0.4°
$S_2-C_2-N_1$	114.5	0.5
$S_2-C_2-N_2$	123.0	0.5
$N_1-C_2-N_2$	122.5	0.6
$C_2-N_1 \cdots O_1$	121.7	
$C_2-N_1 \cdots O_2$	115.2	
$C_2-N_2 \cdots O_1$	116.8	
$C_2-N_2 \cdots O_2$	128.9	
$O_1 \cdots N_1 \cdots O_2$	120.4	
$O_1^1 \cdots N_2 \cdots O_2^4$	111.7	
$O_1-S_1-O_2$	109.8	0.3
$O_1-S_1-O_1^1$	109.2	0.3
$O_1-S_1-O_1^1$	109.9	0.4
$O_2-S_1-O_2^1$	109.0	0.4

In view of the calculated standard deviations S_2 has been given three times the weight of the other atoms. The distances of S_2 , C_2 , N_1 and N_2 from this plane are 0.002, 0.016, 0.005 and 0.005 Å respectively. C_1 lies at 0.56 Å from this plane so that the C_1-S_2 bond makes an angle of 18° with it.

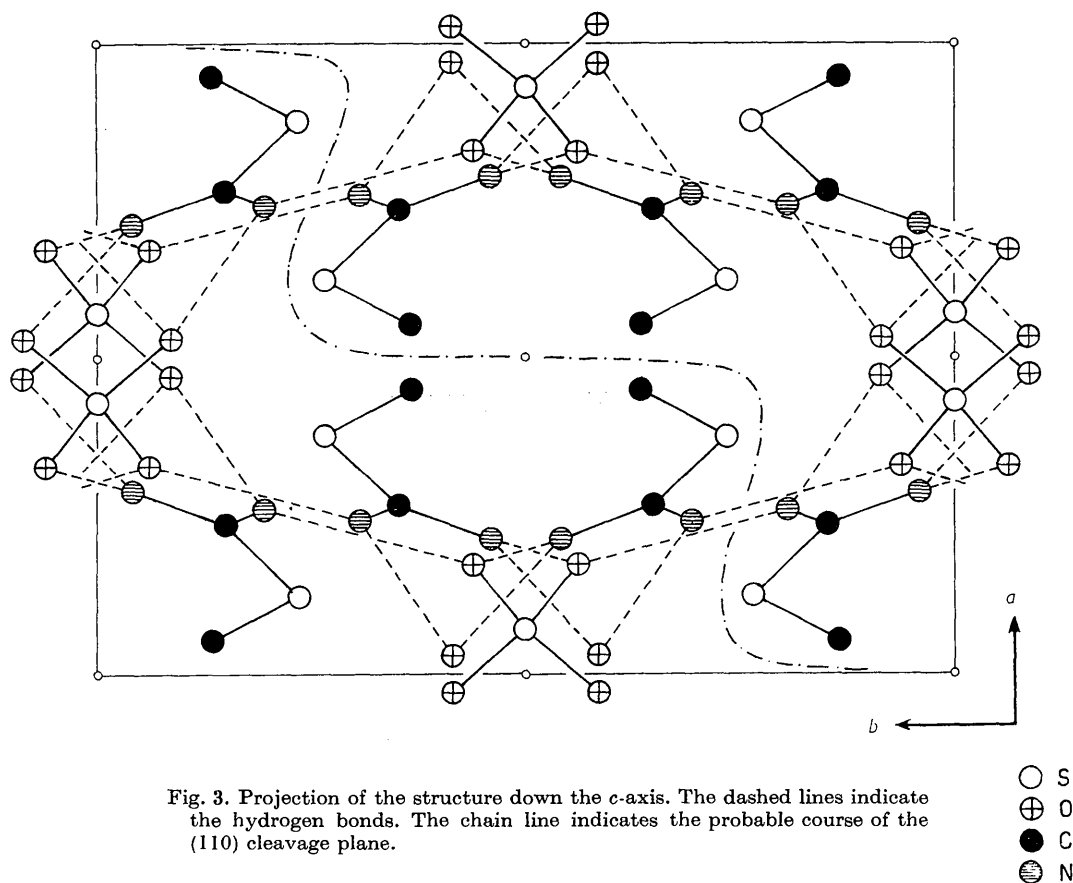


Fig. 3. Projection of the structure down the *c*-axis. The dashed lines indicate the hydrogen bonds. The chain line indicates the probable course of the (110) cleavage plane.

Table 5. *Temperature parameters* (\AA^2)

	U_{11}	U_{22}	U_{33}	$2U_{12}$	$2U_{23}$	$2U_{13}$
S_1	0.0431	0.0260	0.0227	0.0000	-0.0019	0.0000
S_2	0.0710	0.0414	0.0541	0.0224	-0.0013	0.0402
	<i>U</i>					
	O_1	0.047				
	O_2	0.043				
	N_1	0.049				
	N_2	0.041				
	C_1	0.056				
	C_2	0.039				

Table 6. *Direction cosines and mean square amplitudes of vibration for the axes of the vibration ellipsoids of S_1 and S_2*

	$\cos \alpha$	$\cos \beta$	$\cos \gamma$	\bar{u}^2
S_1	1	0	0	0.043
	0	0.966	-0.257	0.026
	0	0.257	0.966	0.022
S_2	0.832	-0.201	-0.517	0.086
	0.276	-0.659	0.700	0.047
	0.481	0.725	0.493	0.034

The C_1 - S_2 bond length of 1.79 \AA is that of a single bond, the single-bond distances for various compounds ranging from 1.78-1.84 \AA (see Abrahams, 1956). The C_2 - S_2 bond of 1.74 \AA is 0.03 \AA longer than the cor-

responding bond in thiourea (Kunchur & Truter, 1958), the difference being possibly significant. As little is known about the nature of the bonds involved detailed discussion of the results is hardly worth while. One may state that the C_2 - S_2 bond has some double-bond character, probably a little less than in thiourea.

The C-S-C angle of 104° agrees very well with those in several other compounds in which sulphur is bound to two carbon atoms (see Sutton, 1958).

The C-N bonds do not differ significantly amongst themselves or from those in thiourea and are of intermediate double-bond character. There is, however, a large unexplained difference between the two S-C-N angles.

The four N-O-hydrogen bonds are medium strong and make angles of 6, 8½, 0 and 14° (in the order of Table 3) with the plane of the thiourea part of the molecule.

The SO_4 ion does not deviate significantly from a regular tetrahedron. The average S-O distance is 1.463 (± 0.004) \AA . This distance is considerably shorter than the S-O distance of 1.488 (± 0.007) \AA found by Atoji & Rundle (1958) in gypsum and that of 1.484 \AA (no accuracy stated) found by Okaya *et al.* (1957) in monomethylammonium aluminium sulphate alum. The difference between our value and that of Atoji

& Rundle is three times the calculated standard deviation for the difference (0.008 Å) and thus would be significant. Other accurate data for the SO₄ ion are those of Singer & Cromer (1959) for zirconium sulphate; they found two significantly different S–O distances of 1.443 and 1.486 Å. It is possible that the calculated standard deviations are too low so that the above differences are not significant. It is also possible that one has to abandon the idea of the SO₄ ion being a regular tetrahedron of constant dimensions under all circumstances. More work on the SO₄ ion is needed to settle this question.

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Search for an Anisotropic Debye-Waller Factor in Cubic Copper-Base Solid Solutions

BY R. W. CAHN

Department of Physical Metallurgy, University of Birmingham, Birmingham 15, England

AND R. FEDER

Pitman-Dunn Laboratories, Frankford Arsenal, Philadelphia 37, Pa., U.S.A.

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A re-examination has been made of the Debye-Waller temperature factor $\exp\{-2B(\sin\theta/\lambda)^2\}$, where B , according to Weiss and co-workers, is a varying function of $\{hkl\}$ when measurements are made on α brass powder. The integrated and peak intensities of several different lines from powder samples of α brass and α aluminium bronze were obtained at 77 °K. and room temperature, with Mo $K\alpha$ radiation, after heat treatments designed to induce varying amounts of short-range order. It was found that deviations of B from isotropy were marginal and not significant except possibly for the 222 lines. No deviations comparable in magnitude with those reported by Weiss and co-workers were found.

Weiss and co-workers (1956) reported measurements of integrated intensities of the 222 and 400 Mo $K\alpha$ lines diffracted by α brass powder at 77 and 295 °K., and showed that these intensities could not be reconciled with the conventional expression $\exp\{-2B(\sin\theta/\lambda)^2\}$ for the Debye-Waller attenuation factor associated with thermal vibrations (hereafter abbreviated DWF), unless B was taken to be a function of direction in the crystal, so that for a given temperature B was different for 222 and 400. They reached this conclusion both by comparing intensities of the same line at two

temperatures, and also by comparing intensities of the two lines at one temperature (and comparing these ratios with computed ratios obtained for various assumed values of B). Values of B were expressed in terms of the Debye characteristic temperature Θ , which accordingly turns out to be effectively a function of crystal direction (we refer to this as an 'anisotropy of the Debye temperature'); values of Θ so obtained do not agree with those calculated from calorimetric measurements on α brass.

Weiss and co-workers believed that the anisotropy